

Coherence lengths for superconductivity in the two-orbital negative- U Hubbard model

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We study the peculiarities of coherency in the superconductivity of two-orbital system. The superconducting phase transition is caused here by the on-site intra-orbital attractions (negative- U Hubbard model) and inter-orbital pair-transfer interaction. The dependencies of critical and non-critical correlation lengths on interaction channels and band fillings are analyzed.

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I. INTRODUCTION

Multi-component superconductivity models, developed more than fifty years starting from the papers [1–3], include rather varied physics. In connection with the presence of interacting order parameters the main properties of the multi-band systems are quite different from the corresponding characteristics in single-band superconductors. The examination of these peculiarities has been an object of growing interest.

The various multi-component theoretical scenarios have been applied for a number of superconducting materials (see Refs. [4–7] and references therein). The research activity in this direction has been especially stimulated by the acceptance of the multi-gap superconductivity in MgB_2 ⁸, cuprates⁹ and iron-arsenic compounds¹⁰. In particular, the derivation of high-quality superconducting regions from oxygen ordering, observed recently in $\text{La}_2\text{CuO}_{4+y}$ ¹¹, supports the multi-band theoretical scenario of superconductivity in cuprates.

In the present contribution we study the coherency of the superconducting ordering of a two-band (two-orbital) system with the negative- U Hubbard intra-orbital pairing and inter-orbital pair-transfer interaction. One can distinguish here two characteristic length scales¹² in the spatial behaviour of superconducting fluctuations. One of these lengths as a function of temperature behaves critically diverging at the phase transition point. The other one remains finite and its temperature dependence is weaker. The formation of these length scales is caused by the interband interaction mixing the superconducting order parameters of initially non-interacting bands. As a result the critical and non-critical coherence lengths associate with critical and non-critical fluctuations (see e.g. Refs. [13,14]) which appear as the certain linear combinations of the deviations from the equilibrium band superconducting orders. Consequently, these length scales cannot be attributed to different bands involved¹⁵ (see also Refs. [16–18]). Our results have been obtained using the superconducting negative- U Hubbard model¹⁹ for a two-orbital system on a two-dimensional lattice.

II. TWO-ORBITAL MODEL OF SUPERCONDUCTIVITY

We start with the Hamiltonian of the two-orbital model²⁰ of the following form:

$$H = \sum_{\alpha} \sum_{i,j} \sum_{\sigma} [t_{ij}^{\alpha\alpha} + (\varepsilon_{\alpha}^0 - \mu) \delta_{ij}] a_{i\alpha\sigma}^{\dagger} a_{j\alpha\sigma} + \frac{1}{2} \sum_{\alpha} \sum_i \sum_{\sigma} U^{\alpha\alpha} n_{i\alpha\sigma} n_{i\alpha-\sigma} + \frac{1}{2} \sum_{\alpha,\alpha'} \sum_i \sum_{\sigma} U^{\alpha\alpha'} a_{i\alpha\sigma}^{\dagger} a_{i\alpha'\sigma} a_{i\alpha-\sigma}^{\dagger} a_{i\alpha'-\sigma}, \quad (1)$$

where $a_{i\alpha\sigma}^{\dagger}$ ($a_{i\alpha\sigma}$) is the electron creation (destruction) operator in the orbital $\alpha = 1, 2$ localized at the site i ; σ is the spin index; $t_{ij}^{\alpha\alpha}$ is the hopping integral; ε_{α}^0 is the orbital energy; μ is the chemical potential; $U^{\alpha\alpha} < 0$ is the intra-orbital attraction energy; $n_{i\alpha\sigma} = a_{i\alpha\sigma}^{\dagger} a_{i\alpha\sigma}$ is the particle number operator; $U^{\alpha\alpha'}$ with $\alpha \neq \alpha'$ is the inter-orbital interaction energy.

The transformation of the Hamiltonian (Eq. 1) into the reciprocal space leaves us with the expression

$$H = \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\sigma} [\varepsilon_{\alpha}(\mathbf{k}) - \mu] a_{\alpha\mathbf{k}\sigma}^{\dagger} a_{\alpha\mathbf{k}\sigma} + \frac{1}{2N} \sum_{\alpha,\alpha'} \sum_{\mathbf{k},\mathbf{k}'} \sum_{\mathbf{q}} \sum_{\sigma} U^{\alpha\alpha'} a_{\alpha(\mathbf{k}+\mathbf{q})\sigma}^{\dagger} a_{\alpha'\mathbf{k}\sigma} \times a_{\alpha(\mathbf{k}'-\mathbf{q})-\sigma}^{\dagger} a_{\alpha'\mathbf{k}'-\sigma}, \quad (2)$$

where $\varepsilon_{\alpha}(\mathbf{k})$ is the electron band energy with a wave vector \mathbf{k} , associated with the orbital α and N is the number of lattice sites (number of atoms).

For the description of spatially homogeneous superconductivity one introduces the equilibrium mean-field Hamiltonian (the terms which do not contain operators have been omitted)

$$H_{mf} = \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\sigma} \tilde{\varepsilon}_{\alpha}(\mathbf{k}) a_{\alpha\mathbf{k}\sigma}^{\dagger} a_{\alpha\mathbf{k}\sigma} + \frac{1}{2} \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\sigma} (\Delta_{\alpha} a_{\alpha\mathbf{k}\sigma}^{\dagger} a_{\alpha-\mathbf{k}-\sigma}^{\dagger} + h.c.) \quad (3)$$

with

$$\tilde{\varepsilon}_\alpha(\mathbf{k}) = \varepsilon_\alpha(\mathbf{k}) + \frac{1}{2}U^{\alpha\alpha}n_\alpha - \mu. \quad (4)$$

Here the average number of electrons per site

$$\sum_\alpha n_\alpha = N^{-1} \sum_\alpha \sum_{\mathbf{k}} \sum_\sigma \langle a_{\alpha\mathbf{k}\sigma}^\dagger a_{\alpha\mathbf{k}\sigma} \rangle_{H_{mf}} \quad (5)$$

determines the position of the chemical potential μ and $\langle \dots \rangle_H = Z^{-1} \text{Sp} \dots \exp(-H/k_B T)$ denotes averaging procedure, where Z is a partition function, k_B is the Boltzmann constant, and T is temperature. The homogeneous equilibrium superconductivity gaps are defined as

$$\Delta_\alpha = N^{-1} \sum_{\alpha'} U^{\alpha\alpha'} \sum_{\mathbf{k}} \langle a_{\alpha'-\mathbf{k}\downarrow} a_{\alpha'\mathbf{k}\uparrow} \rangle_{H_{mf}}. \quad (6)$$

The Eqs. (5) and (6) should be solved self-consistently. It is easy to find that

$$\langle a_{\alpha\mathbf{k}\sigma}^\dagger a_{\alpha\mathbf{k}\sigma} \rangle_{H_{mf}} = \frac{1}{2} \left[1 - \frac{\tilde{\varepsilon}_\alpha(\mathbf{k})}{E_\alpha(\mathbf{k})} \tanh \frac{E_\alpha(\mathbf{k})}{2k_B T} \right] \quad (7)$$

and

$$\langle a_{\alpha-\mathbf{k}\downarrow} a_{\alpha\mathbf{k}\uparrow} \rangle_{H_{mf}} = \frac{-\Delta_\alpha}{2E_\alpha(\mathbf{k})} \tanh \frac{E_\alpha(\mathbf{k})}{2k_B T}, \quad (8)$$

where

$$E_\alpha(\mathbf{k}) = \sqrt{\tilde{\varepsilon}_\alpha^2(\mathbf{k}) + |\Delta_\alpha|^2}, \quad (9)$$

The system of gap equations (6) may be now presented in the form

$$\Delta_\alpha = - \sum_{\alpha'} U^{\alpha\alpha'} \Delta_{\alpha'} \zeta_{\alpha'}(T, \Delta_{\alpha'}) \quad (10)$$

with

$$\zeta_\alpha(T, \Delta_\alpha) = (2N)^{-1} \sum_{\mathbf{k}} \frac{1}{E_\alpha(\mathbf{k})} \tanh \frac{E_\alpha(\mathbf{k})}{2k_B T}, \quad (11)$$

The critical temperature of the phase transition T_c in the two-gap superconductor (Eq. 1) under consideration is determined by the following condition

$$\begin{vmatrix} 1 + U^{11}g_1(T_c) & U^{12}g_2(T_c) \\ U^{21}g_1(T_c) & 1 + U^{22}g_2(T_c) \end{vmatrix} = 0 \quad (12)$$

with $g_\alpha(T_c) = \zeta_\alpha(T_c, 0)$.

The calculated [23] dependency of superconducting phase transition temperature T_c on band filling n is depicted in Fig. 1. As expected, Van Hove singularity present in the middle of the band^{20,21} is reflected as a peak in T_c for half filled system ($n = 1$) where the chemical potential μ passes this singular point. The dependence of the small interband interaction U^{12} is presented by three separate lines $|U^{12}| = |U^{21}| = 0.01, 0.04$, and

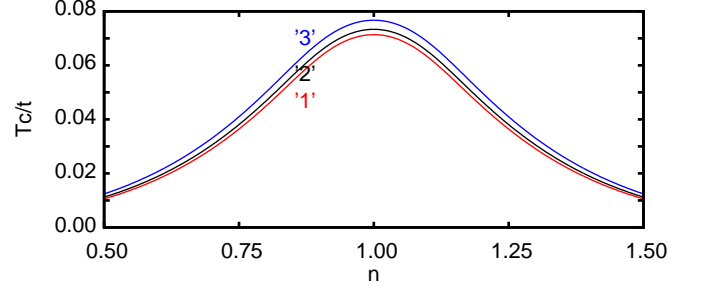


FIG. 1: The critical temperature T_c versus band filling n (equal for each band $n = n_1 = n_2$) for the chosen set of interactions: $U^{11} = -1.2t$, $U^{22} = -1.3t$. $|U^{12}| = |U^{21}|$ was 0.01, 0.04, and 0.07 t for '1'-'3', respectively.

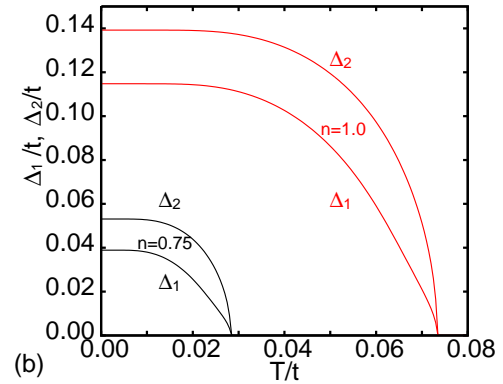
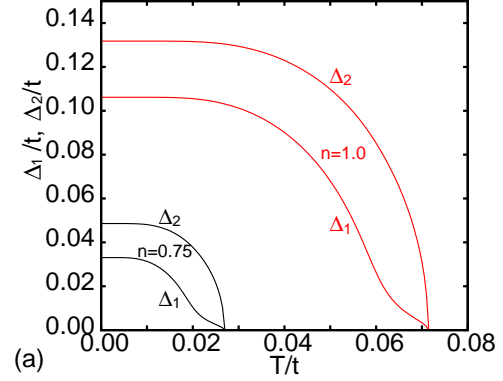


FIG. 2: Superconducting gap parameters Δ_1 and Δ_2 versus temperature T for various band fillings. $U^{11} = -1.2t$, $U^{22} = -1.3t$. $|U^{12}| = |U^{21}| = 0.01, 0.04 t$, for (a) and (b), respectively.

0.07 t for '1'-'3', respectively. The larger $|U^{12}|$ increases the critical temperature T_c .

The superconducting gaps as temperature functions are shown in Fig. 2 for various inter-orbital interactions and band fillings. Note that for fairly low temperature ($T \ll T_c$), Δ_1 and Δ_2 are simultaneously scaled by different intra-orbital interactions U_{ii} ($U_{11} \neq U_{22}$) and the density of states around the chemical potential μ (de-

pendent on band filling which was here $n = 1.0$ or 0.75). On the other hand, the influence of U^{12} is visible in the vicinity of T_c . The stronger U^{12} , the larger Δ_1 . This is a consequence of the "inter-orbital proximity effect".

III. GINZBURG-LANDAU EQUATIONS

In the non-homogeneous situation one has to retain the non-zero momentum of Cooper pairs in the Hamiltonian (Eq. 2) which can be rewritten as

$$\begin{aligned} H = & \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\sigma} [\varepsilon_{\alpha}(\mathbf{k}) - \mu] a_{\alpha\mathbf{k}\sigma}^+ a_{\alpha\mathbf{k}\sigma} \\ & + \frac{1}{2N} \sum_{\alpha, \alpha'} \sum_{\mathbf{k}, \mathbf{k}'} \sum_{\mathbf{q}} \sum_{\sigma} U^{\alpha\alpha'} a_{\alpha\mathbf{k}\sigma}^+ a_{\alpha'(-\mathbf{k}+\mathbf{q})-\sigma}^+ \\ & \times a_{\alpha'(-\mathbf{k}'+\mathbf{q})-\sigma} a_{\alpha\mathbf{k}'\sigma}. \end{aligned} \quad (13)$$

The Bogoliubov theorem determines the upper limit for the free energy of the system

$$F_B = -k_B T \ln(Z'_{mf}) + \langle H - H'_{mf} \rangle_{H'_{mf}}, \quad (14)$$

$$Z'_{mf} = \text{Sp} \exp \left(-\frac{H'_{mf}}{k_B T} \right) \quad (15)$$

by means of the non-equilibrium mean-field Hamiltonian

$$\begin{aligned} H'_{mf} = & \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\sigma} \tilde{\varepsilon}_{\alpha}(\mathbf{k}) a_{\alpha\mathbf{k}\sigma}^+ a_{\alpha\mathbf{k}\sigma} \\ & + \frac{1}{2} \sum_{\alpha} \sum_{\mathbf{k}, \mathbf{q}} \sum_{\sigma} \left(\delta_{\alpha\mathbf{q}} a_{\alpha\mathbf{k}\sigma}^+ a_{\alpha(-\mathbf{k}+\mathbf{q})-\sigma}^+ + h.c. \right) \end{aligned} \quad (16)$$

with non-equilibrium superconductivity gap order parameters $\delta_{\alpha\mathbf{q}}$. The free energy upper limit F_B can be found as an expansion in powers of $\delta_{\alpha\mathbf{q}}$ and \mathbf{q} . After the minimization of F_B and certain transformations one can obtain the system of equations for the non-homogeneous superconductivity gaps $\Delta_{\alpha}(\mathbf{r})$ (the Ginzburg-Landau equations) in the following form:

$$\begin{aligned} \Delta_{\alpha}(\mathbf{r}) = & - \sum_{\alpha'} U^{\alpha\alpha'} \left[g_{\alpha'}(T) - \nu_{\alpha'}(T) |\Delta_{\alpha'}(\mathbf{r})|^2 \right. \\ & \left. + \sum_{l, l'=1}^d \beta_{\alpha' l l'}(T) \frac{\partial}{\partial x_l} \frac{\partial}{\partial x_{l'}} \right] \Delta_{\alpha'}(\mathbf{r}), \end{aligned} \quad (17)$$

where l refers to the Cartesian axis and d is the dimension of lattice. In Eqs. (17)

$$g_{\alpha}(T) = \frac{1}{2N} \sum_{\mathbf{k}} \frac{1}{\tilde{\varepsilon}_{\alpha}(\mathbf{k})} \tanh \frac{\tilde{\varepsilon}_{\alpha}(\mathbf{k})}{2k_B T}, \quad (18)$$

$$\nu_{\alpha}(T) = \frac{-1}{2N} \sum_{\mathbf{k}} \frac{\partial}{\partial |\Delta_{\alpha}|^2} \left[\frac{1}{E_{\alpha}(\mathbf{k})} \tanh \frac{E_{\alpha}(\mathbf{k})}{2k_B T} \right]_{\Delta_{\alpha}=0} \quad (19)$$

$$\begin{aligned} \beta_{\alpha l l'}(T) = & \frac{-1}{4N} \sum_{\mathbf{k}} \frac{\partial}{\partial q_l} \frac{\partial}{\partial q_{l'}} \left\{ \frac{1}{\tilde{\varepsilon}_{\alpha}(\mathbf{k}) + \tilde{\varepsilon}_{\alpha}(\mathbf{k} - \mathbf{q})} \right. \\ & \left. \times \left[\tanh \left(\frac{\tilde{\varepsilon}_{\alpha}(\mathbf{k})}{2k_B T} \right) + \tanh \left(\frac{\tilde{\varepsilon}_{\alpha}(\mathbf{k} - \mathbf{q})}{2k_B T} \right) \right] \right\}_{\mathbf{q}=0}. \end{aligned} \quad (20)$$

By supposing further that

$$\beta_{\alpha l l'} = \beta_{\alpha} \delta_{l l'} \quad (21)$$

one obtains

$$\begin{aligned} \Delta_{\alpha}(\mathbf{r}) = & - \sum_{\alpha'} U^{\alpha\alpha'} \left[g_{\alpha'}(T) - \nu_{\alpha'}(T) |\Delta_{\alpha'}(\mathbf{r})|^2 \right. \\ & \left. + \beta_{\alpha'}(T) \sum_{l=1}^d \frac{\partial^2}{\partial x_l^2} \right] \Delta_{\alpha'}(\mathbf{r}). \end{aligned} \quad (22)$$

IV. COHERENCE LENGTHS

The small deviations from the bulk values of superconductivity gaps

$$\eta_{\alpha}(\mathbf{r}) = \Delta_{\alpha}(\mathbf{r}) - \Delta_{\alpha} \quad (23)$$

satisfy linearized equations

$$\eta_{\alpha}(\mathbf{r}) = - \sum_{\alpha'} U^{\alpha\alpha'} \left[\tilde{g}_{\alpha'}(T) + \beta_{\alpha'}(T) \sum_{l=1}^d \frac{\partial^2}{\partial x_l^2} \right] \eta_{\alpha'}(\mathbf{r}), \quad (24)$$

where

$$\tilde{g}_{\alpha}(T) = g_{\alpha}(T) - 3\nu_{\alpha}(T) (\Delta_{\alpha}(T))^2 \quad (25)$$

with $\Delta_{\alpha}(T)$ being the solutions of the system of equations (Eq. 6). We introduce the coherence lengths ξ as the spatial length scales of $\eta_{\alpha}(\mathbf{r})$:

$$\eta_{1,2}(\mathbf{r}) \sim \exp \left(-\frac{\sum_{l=1}^d x_l}{\sqrt{d}\xi} \right). \quad (26)$$

From Eqs. (22) and (26) follows the system

$$\begin{aligned} & \left(1 + U^{11} \tilde{g}_1(T) + U^{11} \beta_1(T) \xi^{-2} \right) \eta_1(\mathbf{r}) \\ & + \left(U^{12} \tilde{g}_2(T) + U^{12} \beta_2(T) \xi^{-2} \right) \eta_2(\mathbf{r}) = 0 \\ & \left(U^{21} \tilde{g}_1(T) + U^{21} \beta_1(T) \xi^{-2} \right) \eta_1(\mathbf{r}) \\ & + \left(1 + U^{22} \tilde{g}_2(T) + U^{22} \beta_2(T) \xi^{-2} \right) \eta_2(\mathbf{r}) = 0. \end{aligned} \quad (27)$$

the condition for the existence of non-trivial solutions of

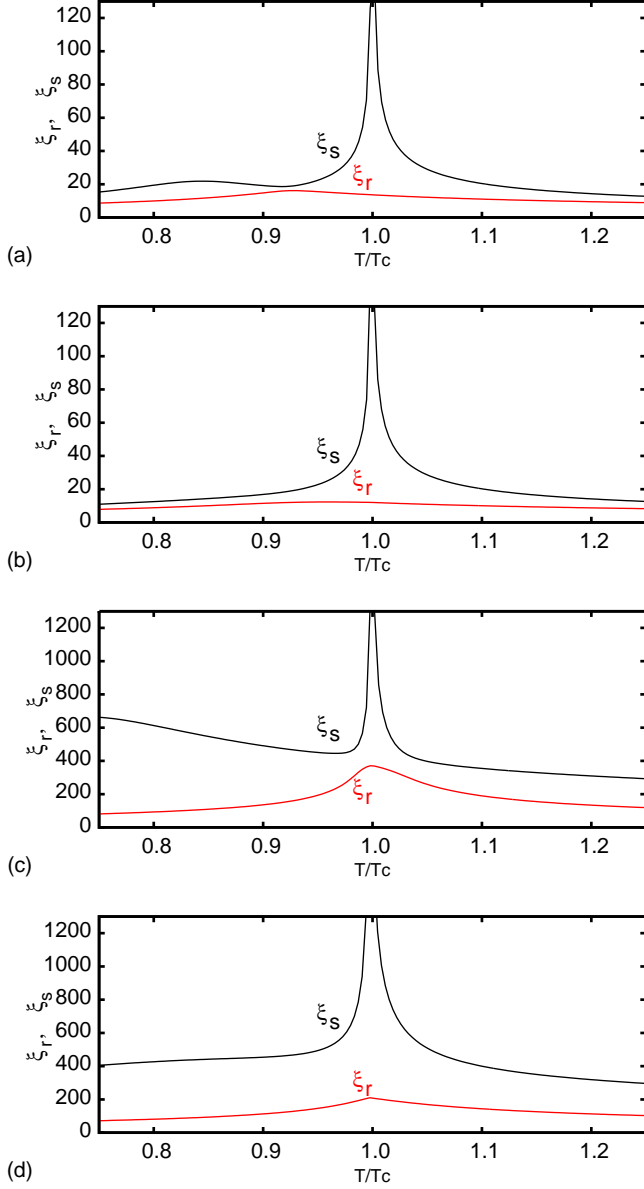


FIG. 3: The soft and rigid coherence lengths (ξ_r , ξ_s) versus temperature T for several sets of parameters: (a) and (c) $|U^{12}| = 0.01$; (b) and (d) $|U^{12}| = 0.04$; (a) and (b) $n = 1.0$, (c) and (d) $n = 0.75$. ξ_r , ξ_s are expressed in terms of lattice constant a .

which,

$$\begin{vmatrix} 1 + U^{11}\tilde{g}_1(T) + U^{11}\frac{\beta_1(T)}{\xi^2}, & U^{12}\tilde{g}_2(T) + U^{12}\frac{\beta_2(T)}{\xi^2} \\ U^{21}\tilde{g}_1(T) + U^{21}\frac{\beta_1(T)}{\xi^2}, & 1 + U^{22}\tilde{g}_2(T) + U^{22}\frac{\beta_2(T)}{\xi^2} \end{vmatrix} = 0, \quad (28)$$

determines the coherence lengths of the two-gap superconductor under consideration. On the basis of Eq. (28) we have two solutions for ξ ,

$$\xi_{s,r}^2(T) = \frac{G(T) \pm \sqrt{G^2(T) - 4K(T)\gamma(T)}}{2K(T)}, \quad (29)$$

where

$$\begin{aligned} G(T) &= (U^{12})^2 [\tilde{g}_1(T)\beta_2(T) + \tilde{g}_2(T)\beta_1(T)] \\ &\quad - [1 + U^{11}\tilde{g}_1(T)] U^{22}\beta_2(T) \\ &\quad - [1 + U^{22}\tilde{g}_2(T)] U^{11}\beta_1(T), \end{aligned} \quad (30)$$

$$\begin{aligned} K(T) &= [1 + U^{11}\tilde{g}_1(T)] [1 + U^{22}\tilde{g}_2(T)] \\ &\quad - (U^{12})^2 \tilde{g}_1(T)\tilde{g}_2(T), \end{aligned} \quad (31)$$

$$\gamma(T) = [U^{11}U^{22} - (U^{12})^2] \beta_1(T)\beta_2(T). \quad (32)$$

The results of our specific two-dimensional lattice calculation of ξ^s and ξ^r versus T/T_c are presented in Fig. 3. The soft or critical coherence length $\xi_s(T)$ diverges at the phase transition point $T = T_c$, while the rigid or non-critical coherence length $\xi_r(T)$ remains finite (see Fig. 3). Furthermore, for fairly small inter-orbital interaction $|U^{12}| = 0.01$ one can notice that ξ_s shows the second maximum (see the peak of ξ_s in Fig. 3a and the rise of ξ_s with lowering T in Fig. 3c) as the memory about the lower transition temperature in the independent orbital (the one with smaller intra-orbital attraction). Such a behaviour is absent for stronger inter-orbital interaction. Additionally, it is worth to notice that the coherence length is much larger for $n = 0.75$ than for $n = 1.0$ (see Figs. 3a,b and 3c,d respectively). This is again the effect of the density of states. Evidently, Van Hove singularity²¹ scales strongly T_c and other superconducting parameters including the coherence length.

V. CONCLUSIONS

We considered two orbital superconductor with a small interaction coupling producing possibility of the Copper pair transfer between orbitals. The system have been modelled by the negative- U Hubbard model on the two dimensional lattice. The obtained two values of the coherence lengths reflects the multi-orbital mechanism of superconductivity and could have some consequence in specific spatial properties of such a complex supeconducting state^{15,22}

Our results show that band filling strongly effects on the coherence lengths. Especially, Van Hove singularity present in the density of states plays an important role in scaling the coherence lengths. The temperature dependence of the coherence lengths in our lattice model confirms the similar calculations in the continuous system¹². One solution diverges near the phase transition point, while the other one is non-critical. The non-monotonic temperature dependence of coherence lengths is more pronounced if the inter-orbital coupling is sufficiently weak.

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- ¹ H. Suhl, B. T. Matthias and L. P. Walker, Phys. Rev. Lett. **3**, 552 (1959).
 - ² V. A. Moskalenko, Fiz. Met. Metalloved. **8**, 503 (1959).
 - ³ J. Kondo, Progr. Theor. Phys. **29**, 1 (1963).
 - ⁴ V. Z. Kresin, S. A. Wolf, Phys. Rev. B **46**, 6458 (1992).
 - ⁵ N. Kristoffel, P. Konsin, T. Örd, Riv. Nuovo Cimento **19**(8), 1 (1994).
 - ⁶ A. Bianconi, J. Superconductivity **18**, 625 (2005).
 - ⁷ R. Caivano, M. Fratini, N. Poccia, A. Ricci, A. Puri, Z.-A. Ren, X.-L. Dong, J. Yang, W. Lu, Z.-X. Zhao, L. Barba, A. Bianconi, Supercond. Sci. Technol. **22**, 014004 (2009).
 - ⁸ S. Tsuda, T. Yokoya, Y. Takano, H. Kito, A. Matsushita, F. Yin, J. Itoh, H. Harima, S. Shin, Phys. Rev. Lett. **91**, 127001 (2003).
 - ⁹ R. Khasanov, A. Shengelaya, A. Maisuradze, F. La Matina, A. Bussmann-Holder, H. Keller, K. A. Müller, Phys. Rev. Lett. **98**, 057007 (2007).
 - ¹⁰ H. Ding, P. Richard, K. Nakayama, K. Sugawara, T. Arakane, Y. Sekiba, A. Takayama, S. Souma, T. Sato, T. Takahashi, Z. Wang, X. Dai, Z. Fang, G. F. Chen, J. L. Luo, N. L. Wang, EPL **83**, 47001 (2008).
 - ¹¹ N. Poccia, M. Fratini, A. Ricci, G. Campi, L. Barba, A. Vittorini-Orgeas, G. Bianconi, G. Aeppli, A. Bianconi, Nature Materials **10**, 733 (2011).
 - ¹² T. Örd, K. Rågo, A. Vargunin, J. Supercond. Novel Magn., accepted for publication.
 - ¹³ T. Örd, K. Rågo, A. Vargunin, J. Supercond. Novel Magn., **22**, 85 (2009).
 - ¹⁴ T. Örd, K. Rågo, A. Vargunin, In: Physical Properties of Nanosystems, Ed. J. Bonča, S. Kruchinin, NATO Science for Peace and Security Series B: Physics and Biophysics, p.177, Springer, Dordrecht (2011).
 - ¹⁵ J. Carlström, E. Babaev, M. Speight, Phys. Rev. B **83**, 174509 (2011).
 - ¹⁶ Y. M. Poluektov, V. V. Krasilnikov, Fizika Nizkhik Temperatur **15**, 1251 (1989).
 - ¹⁷ P. Konsin, phys. stat. sol. (b) **189**, 185 (1995).
 - ¹⁸ N. Kristoffel, T. Örd, P. Rubin, Supercond. Sci. Technol., **22**, 014006 (2009).
 - ¹⁹ R. Micnas, J. Ranninger, and S. Robaszkiewicz, Rev. Mod. Phys. **62**, 113 (1990).
 - ²⁰ A. Ciechan, K.I. Wysokiński, Phys. Rev. B **80**, 224523 (2009).
 - ²¹ R. S. Markiewicz, J. Phys. Chem. Solids **58** 1179 (1997).
 - ²² J.F. Annett, G. Litak, B.L. Gyorffy, K.I. Wysokiński, Phys. Rev. B **66**, 134514 (2002).
 - ²³ All illustrative numerical calculations have been carried out for two-dimensional square lattice with hopping integrals between nearest neighbours $t_{ij}^{\alpha\alpha} = t$. The electron band energies associated with s -orbitals are $\varepsilon_1(\mathbf{k}) = \varepsilon_2(\mathbf{k}) = \varepsilon^0 - 2t [\cos(ak_x) + \cos(ak_y)]$ with $-\frac{\pi}{a} \leq k_{x,y} \leq \frac{\pi}{a}$ where a is the lattice constant. We have neglected the Hartree-type term in Eq. (4).